

SAMPLE 18.15 Kinematics of a four bar linkage. A four bar linkage ABCD is shown in the figure (fourth bar is the ground AD) at some instant t_0 . The driving bar AB rotates with angular velocity $\bar{\omega}_{AB} = \dot{\theta}(t)\hat{k}$. Find the angular velocities of rods BC and CD as a function of $\dot{\theta}$. How can you solve for the positions of the bars at any t if the initial configuration is as shown in the figure?

Solution Let the angles that rods AB, BC, and CD make with the horizontal (x -axis) be θ , β , and ϕ , respectively. Then, we can write $\bar{\omega}_{BC} = \dot{\beta}\hat{k}$ and $\bar{\omega}_{CD} = \dot{\phi}\hat{k}$. We have to find $\dot{\beta}$ and $\dot{\phi}$.

Note that the motion of point B is a simple circular motion about point A with given angular velocity $\bar{\omega}_{AB}$. Thus, the velocity of point B is known. Now, we can find the velocity of point C two ways: (i) by considering rod CB: $\bar{v}_C = \bar{v}_B + \bar{v}_{C/B} = \bar{v}_B + \bar{\omega}_{BC} \times \bar{r}_{C/B}$, and (ii) by considering rod CD: $\bar{v}_C = \bar{\omega}_{CD} \times \bar{r}_{C/D}$. Either way the velocity must be the same. Thus, we have a 2-D vector equation with two unknowns $\dot{\beta}$ and $\dot{\phi}$. We can get two independent scalar equations from the vector equation and thus we can solve for the desired unknowns.

Let us use the rotating base vectors $(\hat{\lambda}_1, \hat{n}_1)$, $(\hat{\lambda}_2, \hat{n}_2)$, and $(\hat{\lambda}_3, \hat{n}_3)$ with rods AB, BC, and CD, respectively. Note that these base vectors are basically the $(\hat{e}_r, \hat{e}_\theta)$ pairs; we use $(\hat{\lambda}, \hat{n})$ just for the sake of easy subscripting. Now,

$$\begin{aligned}\bar{v}_B &= \bar{\omega}_{AB} \times \bar{r}_{B/A} = \ell_1 \dot{\theta} \hat{n}_1 \\ \bar{v}_C &= \bar{v}_B + \bar{\omega}_{BC} \times \bar{r}_{C/B} = \ell_1 \dot{\theta} \hat{n}_1 + \ell_2 \dot{\beta} \hat{n}_2\end{aligned}\quad (18.67)$$

$$\text{also, } \bar{v}_C = \bar{\omega}_{CD} \times \bar{r}_{C/D} = \ell_3 \dot{\phi} \hat{n}_3 \quad (18.68)$$

Thus, from eqn. (18.67) and (18.68), we have,

$$\ell_1 \dot{\theta} \hat{n}_1 + \ell_2 \dot{\beta} \hat{n}_2 = \ell_3 \dot{\phi} \hat{n}_3 \quad (18.69)$$

Dotting eqn. (18.69) with $\hat{\lambda}_2$ (to eliminate $\dot{\beta}$ term), we get

$$\begin{aligned}\ell_1 \dot{\theta} (\hat{n}_1 \cdot \hat{\lambda}_2) &= \ell_3 \dot{\phi} (\hat{n}_3 \cdot \hat{\lambda}_2) \\ \Rightarrow \dot{\phi} &= \frac{\ell_1 (\hat{n}_1 \cdot \hat{\lambda}_2)}{\ell_3 (\hat{n}_3 \cdot \hat{\lambda}_2)} \dot{\theta}.\end{aligned}\quad (18.70)$$

Similarly, dotting eqn. (18.69) with $\hat{\lambda}_3$ (to eliminate $\dot{\phi}$ term), we get

$$\dot{\beta} = -\frac{\ell_1 (\hat{n}_1 \cdot \hat{\lambda}_3)}{\ell_2 (\hat{n}_2 \cdot \hat{\lambda}_3)} \dot{\theta}.\quad (18.71)$$

We are practically done at this point with the kinematics — we have found $\dot{\beta}$ and $\dot{\phi}$ as functions of $\dot{\theta}$. The various dot products are just geometry and vector algebra. To write them explicitly, we note that $\hat{\lambda}_1 = \cos \theta \hat{i} + \sin \theta \hat{j}$, $\hat{n}_1 = -\sin \theta \hat{i} + \cos \theta \hat{j}$, $\hat{\lambda}_2 = \cos \beta \hat{i} + \sin \beta \hat{j}$, etc. Thus,

$$\begin{aligned}\hat{n}_1 \cdot \hat{\lambda}_2 &= -\sin \theta \cos \beta + \cos \theta \sin \beta = \sin(\beta - \theta) \\ \hat{n}_3 \cdot \hat{\lambda}_2 &= \sin(\beta - \phi) \\ \hat{n}_1 \cdot \hat{\lambda}_3 &= \sin(\phi - \theta), \quad \hat{n}_2 \cdot \hat{\lambda}_3 = \sin(\phi - \beta).\end{aligned}$$

Substituting the appropriate expressions in eqn. (18.71) and (18.70), we get

$$\dot{\beta} = -\frac{\ell_1 \sin(\phi - \theta)}{\ell_2 \sin(\phi - \beta)} \dot{\theta}, \quad \dot{\phi} = \frac{\ell_1 \sin(\beta - \theta)}{\ell_3 \sin(\beta - \phi)} \dot{\theta}.$$

$$\bar{\omega}_{BC} = -\frac{\ell_1 \sin(\phi - \theta)}{\ell_2 \sin(\phi - \beta)} \dot{\theta} \hat{k}, \quad \bar{\omega}_{CD} = \frac{\ell_1 \sin(\beta - \theta)}{\ell_3 \sin(\beta - \phi)} \dot{\theta} \hat{k}$$

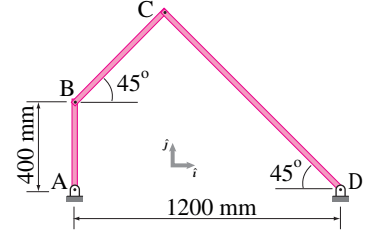


Figure 18.64

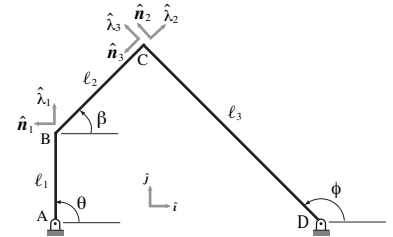


Figure 18.65

Note that the expressions for $\dot{\beta}$ and $\dot{\phi}$ are coupled, nonlinear, first order ordinary differential equations. To be able to find $\dot{\theta}$, $\beta(t)$ and $\phi(t)$, we need to integrate $\dot{\theta}$, $\dot{\beta}$, and $\dot{\phi}$. Here, we set up these differential equations for numerical integration. Although, we can use any given $\dot{\theta}(t)$ (e.g., αt or $\dot{\theta}_0 \sin(\Omega t)$ or whatever), for definiteness in our numerical integration, let us take a constant $\dot{\theta}$, that is, let $\dot{\theta} = C = 10$ rad/s (say). So, our equations are,

$$\dot{\theta} = C, \quad \dot{\beta} = \frac{\ell_1 \sin(\phi - \theta)}{\ell_2 \sin(\beta - \phi)} C, \quad \dot{\phi} = \frac{\ell_1 \sin(\beta - \theta)}{\ell_3 \sin(\beta - \phi)} C,$$

and the initial conditions are $\theta(0) = \pi/2$, $\beta(0) = \pi/4$, $\phi(0) = 3\pi/4$.

Here is a pseudocode that we use to integrate these equations numerically for a period of $2\pi/C = \pi/5$ seconds (one complete revolution of AB).

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ODEs = {thetadot = C,
        betadot = (l1/l2)*sin(phi-theta)/sin(beta-phi)*C,
        phidot = (l1/l3)*sin(beta-theta)/sin(beta-phi)*C}
IC = {theta(0) = pi/2, beta(0) = pi/4, phi(0) = 3*pi/4}
Set C=10, l1=.4, l2=.4*sqrt(2), l3=.8*sqrt(2)
Solve ODEs with IC for t=0 to t=pi/5
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After we get the angles, we can compute the xy coordinates of points B and C at each instant and plot the mechanism at those instants. Plots thus obtained from our numerical solution are shown in fig. 18.66 where the configuration of the mechanism is shown at 9 equally spaced times between $t = 0$ to $t = \pi/5$.

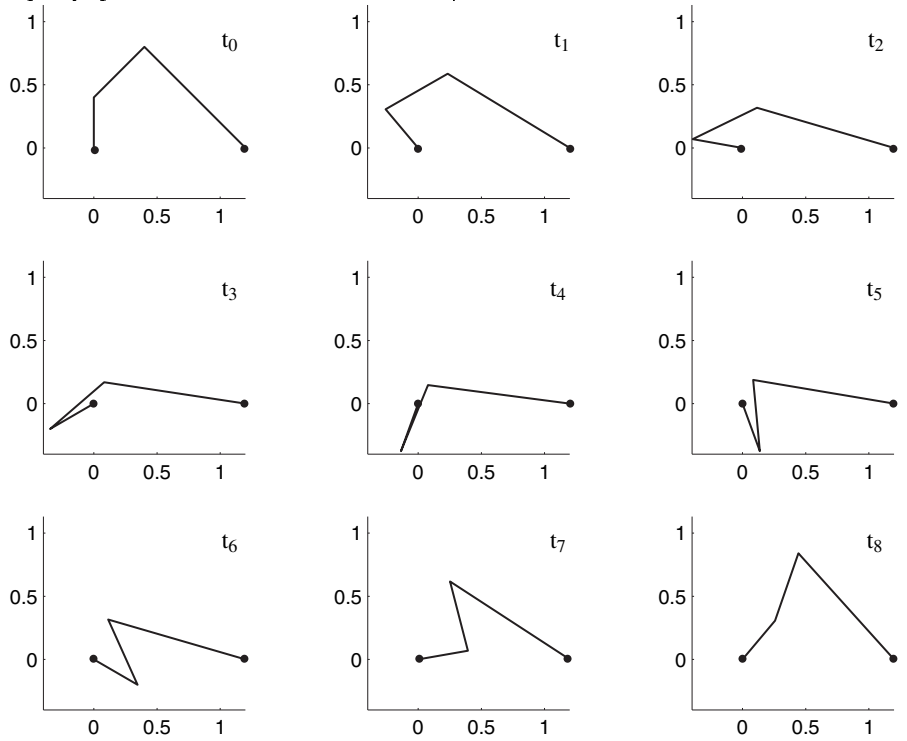


Figure 18.66: Several configurations of the mechanism at equal intervals of time during one complete revolution of the driving link. After t_8 the mechanism returns to the initial configuration.